

TUTORIAL NOTE III

Suppose u is harmonic in Ω . Then integrating by parts we have

$$\int_{\Omega} u \Delta \varphi = 0 \quad \text{for any } \varphi \in C_c^2(\Omega).$$

The converse is also true and due to Wely.

Exercise 1 (Wely's Lemma). *Let $\Omega \subset \mathbb{R}^n$ with $n \geq 3$. Suppose $u \in C(\Omega)$ satisfies*

$$\int_{\Omega} u \Delta \varphi = 0 \quad \text{for any } \varphi \in C_c^2(\Omega). \quad (0.1)$$

Then u is harmonic in Ω .

Proof. First, we claim that for any $B_r(x) \subset \Omega$ there holds

$$r \int_{\partial B_r(x)} u(y) dS_y = n \int_{B_r(x)} u(y) dy. \quad (0.2)$$

Indeed, for simplicity we assume that $x = 0$. Set

$$\varphi_k(y, r) = \begin{cases} (|y|^2 - r^2)^k, & |y| \leq r, \\ 0, & |y| > r, \end{cases}$$

and

$$\psi_k(y, r) = (|y|^2 - r^2)^{k-2} (2(k-1)|y|^2 + n(|y|^2 - r^2))$$

for $k \geq 2$. By an elementary computation, we have $\varphi_k(\cdot, r) \in C_0^2(\Omega)$ and

$$\Delta_y \varphi_k(y, r) = \begin{cases} 2k\psi_k(y, r), & |y| \leq r, \\ 0, & |y| > r. \end{cases}$$

for $k \geq 3$. For $|y| < r$ and $k \geq 3$, we have

$$\begin{aligned} 2k\partial_r \psi_k &= \partial_r \Delta_y \varphi_k \\ &= \Delta_y \partial_r \varphi_k \\ &= -2kr \Delta_y \varphi_{k-1} \\ &= -4k(k-1)r\psi_{k-1}. \end{aligned} \quad (0.3)$$

From (0.1), we see that

$$6 \int_{B_r(0)} u(y) \psi_3(y, r) dy = \int_{B_r(0)} u(y) \Delta_y \varphi_3(y, r) dy = 0.$$

Now we prove

$$\int_{B_r(0)} u(y) \psi_3(y, r) dy = 0 \Rightarrow \int_{B_r(0)} u(y) \psi_2(y, r) dy = 0. \quad (0.4)$$

In fact, we differentiate the first identity in (0.4) with respect to r and get

$$\int_{\partial B_r(0)} u(y) \psi_3(y, r) dy + \int_{B_r(0)} u(y) \partial_r \psi_3(y, r) dy = 0.$$

Note that $\psi_3(y, r) = 0$ for $|y| = r$. Therefore, we have

$$\int_{B_r(0)} u(y) \partial_r \psi_3(y, r) dy = 0.$$

By (0.3), we see that (0.4) is true. We have

$$\int_{B_r(0)} u(y) ((n+2)|y|^2 - nr^2) dy = 0.$$

Differentiating with respect to r again, we obtain (0.2). Note that, from (0.2), we know that

$$\begin{aligned} & \frac{d}{dr} \left(\frac{1}{\omega_n r^{n-1}} \int_{\partial B_r(x)} u(y) dS_y \right) \\ &= \frac{n}{\omega_n} \frac{d}{dr} \left(\frac{1}{r^n} \int_{B_r(x)} u(y) dy \right) \\ &= \frac{n}{\omega_n r^{n+1}} \left(r \int_{\partial B_r(x)} u(y) dS_y - n \int_{B_r(x)} u(y) dy \right) = 0. \end{aligned}$$

This implies

$$u(x) = \lim_{\rho \rightarrow 0} \frac{1}{\omega_n \rho^{n-1}} \int_{\partial B_\rho(x)} u(y) dS_y = \frac{1}{\omega_n r^{n-1}} \int_{\partial B_r(x)} u(y) dS_y,$$

which means that u is harmonic. □

Second, we introduce the removable singularity for Harmonic function.

Exercise 2. Suppose u is harmonic in $B_R(0) \setminus \{0\}$ and satisfies

$$u(x) = \begin{cases} o(\log|x|), & n = 2, \\ o(|x|^{2-n}), & n \geq 3, \end{cases} \quad \text{as } |x| \rightarrow 0.$$

Then u can be defined at 0 so that it is C^2 and harmonic in $B_R(0)$.

Proof. We prove the case $n = 3$ as an example. Assume u is continuous in $0 < |x| \leq R$. Let v solve

$$\begin{cases} \Delta v = 0 & \text{in } B_R, \\ v = u & \text{on } \partial B_R. \end{cases}$$

We claim that $u = v$ in $B_R \setminus \{0\}$. Indeed, we can consider $w = v - u$ in $B_R(0) \setminus \{0\}$ and $M_r = \max_{\partial B_r(0)} |w|$. We observe that

$$|w(x)| \leq M_r r |x|^{-1} \quad \text{on } \partial B_r.$$

Note that w and $|x|^{-1}$ are harmonic in $B_R \setminus B_r$. Hence the maximum principle implies

$$|w(x)| \leq M_r r |x|^{-1} \quad \text{for any } x \in B_R \setminus B_r.$$

Note also that, for all $r \in (0, R)$,

$$M_r \leq \max_{\partial B_r(0)} |u(x)| + \max_{\partial B_r(0)} |v(x)| \leq \max_{\partial B_r(0)} |u(x)| + \max_{\partial B_R(0)} |u(x)|.$$

Combining the above estimates, we have for each fixed $x \neq 0$,

$$|w(x)| \leq |x|^{-1} r \left(\max_{\partial B_r(0)} |u(x)| + \max_{\partial B_R(0)} |u(x)| \right) \rightarrow 0 \text{ as } r \rightarrow 0,$$

that is $w = 0$ in $B_R(0) \setminus \{0\}$. □